Downloaded from UvA-DARE, the institutional repository of the University of Amsterdam (UvA) http://hdl.handle.net/11245/2.46397

File ID uvapub:46397 Filename 210957y.pdf Version unknown

SOURCE (OR PART OF THE FOLLOWING SOURCE):

Type article

Title Effect of Long-Range Potential Fluctuations on Scaling in the Integer

Quantum Hall-Effect

Author(s) H.P. Wei, S.I. Lin, D.C. Tsui, A.M.M. Pruisken

Faculty UvA: Universiteitsbibliotheek

Year 1992

FULL BIBLIOGRAPHIC DETAILS:

http://hdl.handle.net/11245/1.426642

Copyright

It is not permitted to download or to forward/distribute the text or part of it without the consent of the author(s) and/or copyright holder(s), other than for strictly personal, individual use, unless the work is under an open content licence (like Creative Commons).

PHYSICAL REVIEW B

Effect of long-range potential fluctuations on scaling in the integer quantum Hall effect

H. P. Wei,* S. Y. Lin, and D. C. Tsui Department of Electrical Engineering, Princeton University, Princeton, New Jersey 08544

A. M. M. Pruisken

Institute for Theoretical Physics, University of Amsterdam, The Netherlands (Received 1 July 1991; revised manuscript received 5 December 1991)

We report a set of transport data taken in two low-mobility $GaAs/Al_xGa_{1-x}As$ heterostructures. When T > 200 mK, we find that the T dependence of $(d\rho_{xy}/dB)^{max}$ behaves differently in different Landau levels, whereas when T < 200 mK, it behaves like $T^{-0.42}$ as reported by Wei et al. [Phys. Rev. Lett. 61, 1294 (1988)]. The characteristic T(=200 mK) for observing the critical behavior is much lower than that of previous observations in the $In_xGa_{1-x}As/InP$ heterostructure. This lowering of T for scaling is attributed to the dominance of long-range potential fluctuations due to the remote ionized impurities in the $Al_xGa_{1-x}As$.

The appearance of the integer quantum Hall effect (IQHE) in a two-dimensional electron gas (2DEG) in high magnetic fields (B) at low temperatures (T) indicates that the electronic states are localized except at some singular energies, where they are delocalized. The transition of the electronic states at the Fermi level (E_F) from localized to delocalized states, when B is swept through adjacent quantum Hall plateaus, has been found to exhibit scaling behavior.^{2,3} In particular, in an $In_xGa_{1-x}As/InP$ heterostructure, the T dependence of the peak of the derivative of the Hall resistance ρ_{xy} with respect to B $[(d\rho_{xy}/dB)^{\text{max}}]$ diverges like $T^{-\kappa}$ ($\kappa = 0.42$), independent of Landau levels, when T < 4.2 K. This result is a direct consequence of the scaling theory of Pruisken. 4 Basically, the scaling behavior is due to the quantum interference of electrons in a disordered medium. At finite T, the inelastic scattering length (l_{in}) is the largest length scale within which the quantum interference makes sense, and therefore is the effective sample size.⁵ The scaling theory of the IQHE (Refs. 6 and 7) assumes the existence of uncorrelated, δ -function-like potential fluctuations. This is realized in $In_xGa_{1-x}As/InP$ heterostructures where the correlation length of the random potential is approximately equal to the lattice constant. In this paper, we discuss the situation in $GaAs/Al_xGa_{1-x}As$ heterostructures in which the potential fluctuations are long range.

In the GaAs/Al_xGa_{1-x}As samples that we have studied⁸ (all grown by molecular-beam epitaxy), we find that in most cases, the T-dependent behavior of $(d\rho_{xy}/dB)^{\text{max}}$ is sample and Landau-level dependent. A similar result is also reported in a recent paper by Koch $et~al.^9$ On the other hand, in two of our samples, we did find evidence that the $T^{-0.42}$ critical behavior is observed for T < 200 mK. The first sample has an electron density $n = 1.9 \times 10^{11}$ cm⁻² and a mobility $\mu = 55\,000$ cm²/Vs at 4.2 K. Figure 1 shows the transport coefficients ρ_{xx} and ρ_{xy} at 66 mK. There is no signature of the fractional quantum Hall effect that is observed in samples with higher mobility, ¹⁰ thus the many-body effect is not relevant in this sample. The other sample has $n = 2.1 \times 10^{11}$ cm⁻² and $\mu = 65\,000$

cm²/V s at 4.2 K. We find that the transport properties in these two samples are the same in all respects. Therefore, we will present the data from the first sample only.

In Fig. 2 we plot the T dependent $(d\rho_{xy}/dB)^{max}$ for 30 mK < T < 4.2 K in three Landau levels. There are four different symbols in each curve representing data taken from four different runs. We notice that for T > 200 mK, $(d\rho_{xy}/dB)^{max}$ behave differently in each Landau level. However, they behave more or less the same for T < 200 mK. The straight line on top of this figure is drawn for reference purpose. It has a slope of 0.42 on this loglog plot. The data suggest that for T < -200 mK, $(d\rho_{xy}/dB)^{max}$ has a power-law dependence, $T^{-\kappa}$, independent of Landau levels. The exponent $\kappa = 0.42$, the same as that reported previously in the $In_xGa_{1-x}As/InP$ sample. ²

The main difference from the previous $In_xGa_{1-x}As/$ InP results is the value of $T_{\rm sc}$, which is the T where the $T^{-0.42}$ power-law behavior starts to appear $(T_{\rm sc} \sim 200$ mK in the present sample and ~4 K in the $In_xGa_{1-x}As/InP$ sample of Ref. 2). We find that T_{sc} in the $GaAs/Al_xGa_{1-x}As$ sample is substantially lower than T, where the conductivity peak $(\sigma_{xx}^{\text{max}})$ has a maximum (see Fig. 3) and which equals about 1 K. We emphasize that the latter T is solely an effect of the thermal width of the Fermi-Dirac distribution, 2,3 and in $In_xGa_{1-x}As/InP$ it equals $T_{\rm sc}$. Experimentally, transport coefficients are obtained by measuring a voltage across two contacts which are typically a millimeter apart. This procedure is viewed as an average over the macroscopic sample composed of a large number of microscopic effective samples, each with a size of l_{in} . The observation of much lower T_{sc} in the $GaAs/Al_xGa_{1-x}As$ sample is an indication of the presence of long-range potential fluctuations. More specifically, the lack of the power-law behavior in the range of T from 1 K to 200 mK is mainly a result of l_{in} not being much larger than the distance over which the potential fluctuations are correlated.

The qualitative difference in the potential distributions in the $In_xGa_{1-x}As/InP$ and the $GaAs/Al_xGa_{1-x}As$ heterostructures is also obvious from the metallurgy of the

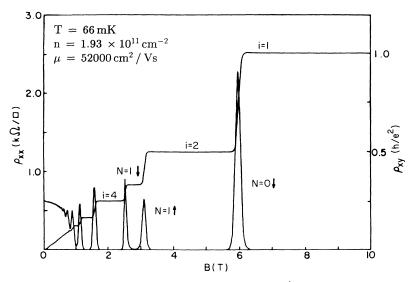


FIG. 1. Magneto-transport coefficients ρ_{xx} and ρ_{xy} vs B at T = 66 mK in a GaAs/Al_xGa_{1-x}As heterostructure. $N = 1 \downarrow$, $1 \uparrow$, and $0 \downarrow$ are the indexes for spin (\downarrow,\uparrow) split Landau levels, in which the E_F resides when B is swept through adjacent quantum Hall plateaus.

sample. Since the 2DEG in the $In_xGa_{1-x}As/InP$ heterostructure is in the $In_xGa_{1-x}As$ layer, which is an alloy, ¹¹ the potential fluctuations are therefore short ranged compared to the cyclotron radius (typically 100 Å) and hence to l_{in} . On the other hand, the 2DEG in the GaAs/ $Al_xGa_{1-x}As$ heterostructures is in the GaAs layer, and it is well known that the dominant scattering mechanism at low T is the remote ionized impurities away from the 2DEG layer. ^{12,13} One should then expect smooth, longrange potential fluctuations. ¹⁴ The dominance of long-

(stimulation of the control of the c

FIG. 2. T dependence of $(d\rho_{xy}/dB)^{\text{max}}$ for three Landau levels, $N=1\downarrow$, $1\uparrow$, and $0\downarrow$. The solid line is drawn for reference purpose and has a slope of 0.42.

range potential fluctuations lowers the T, below which scaling starts and complicates the observability of the critical phenomenon. Given the T available to the experiment, our data is quite limited compared to that of our previous work on $\text{In}_x\text{Ga}_{1-x}\text{As/InP}$, in which case the critical behavior in $d\rho_{xy}/dB$, $(\Delta B)^{-1}$, $d^2\rho_{xx}/dB^2$, and $d^3\rho_{xy}/dB^3$ are well established in more than two decades in T^3

Another striking difference may be found in the spin (\downarrow,\uparrow) dependence of σ_{xx}^{max} corresponding to the $N=1\uparrow$ and $1\downarrow$ Landau levels for T less than ~ 1 K in Fig. 3. They are not equal as in the case of the $\ln_x \text{Ga}_{1-x} \text{As/InP}$ sample. Although there is a clear tendency toward the $T^{-0.42}$ critical behavior in $(d\rho_{xy}/dB)^{\text{max}}$ at low T, the ori-

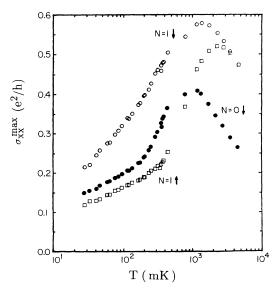


FIG. 3. σ_{xx}^{max} as a function of T for three Landau levels, $N=1\downarrow,1\uparrow$, and $0\downarrow$.

gin of this effect, not relevant to scaling, is not understood at present.

In summary, we may say that the critical behavior in the metal-insulator transition in the IQHE is severely affected by the long-range nature of the potential fluctuations in the $GaAs/Al_xGa_{1-x}As$ samples, in contrast to the short-range random alloy potential in the $In_xGa_{1-x}As/InP$ samples. The dominance of long-range potential fluctuations lowers the T below which the scaling becomes observable. In two of our samples, scaling and the related critical behavior are observed for T less than 200 mK. We conclude from this observation that the apparent lack of scaling in most of our $GaAs/Al_xGa_{1-x}As$ samples is due to crossover effects, which are dependent on the microscopic details of the sample, and much lower T will be needed to observe scaling in these samples.

Note added. While this paper was being reviewed, we

became aware of the recent work by Koch et al. [Phys. Rev. Lett. 67, 833 (1991)], who found that the temperature exponent p for the inelastic scattering length $l_{\rm in} (=T^{-p/2})$ depends on the quality of samples. Thus, they concluded that $\kappa(\sim p)$ is material dependent in the T range available in the laboratory. On the other hand, if $l_{\rm in}$ is due to electron-electron interaction at low T, it is likely that p should not depend on the details of disorder in the limit of very large sample size.

We are indebted to Dr. K. Alavi for the GaAs/ $Al_xGa_{1-x}As$ samples. This work is supported by the Office of Naval Research under Contract No. N00014-89-J-1567. One of us (A.M.M.P) acknowledges the support in part by the Dutch Science Foundation (FOM) and by NSF Grant No. DMR-8907580.

^{*}Present address: Department of Physics, University of Maryland, College Park, MD 20742.

¹M. A. Paalanen, D. C. Tsui, and A. C. Gossard, Phys. Rev. B 25, 5566 (1982).

²H. P. Wei, D. C. Tsui, M. Paalanen, and A. M. M. Pruisken, Phys. Rev. Lett. **61**, 1294 (1988).

³H. P. Wei, S. W. Hwang, D. C. Tsui, and A. M. M. Pruisken, Surf. Sci. **229**, 30 (1990).

⁴A. M. M. Pruisken, Phys. Rev. Lett. 61, 1297 (1988).

⁵D. J. Thouless, Phys. Rev. Lett. 39, 1167 (1977).

⁶A. M. M. Pruisken, Phys. Rev. B 32, 2636 (1985), and references therein.

⁷A. M. M. Pruisken, in *The Quantum Hall Effect*, edited by R. Prange and S. Girvin, Springer Physics Lecture Series

⁽Springer-Verlag, Berlin, 1986).

⁸H. P. Wei, S. Lin, D. C. Tsui, and K. Alavi, Bull. Am. Phys. Soc. 33, 665 (1988); H. Z. Zheng, H. P. Wei, and D. C. Tsui, *ibid.* 32, 892 (1987).

⁹S. Koch, R. J. Haug, K. v. Klitzing, and K. Ploog, Phys. Rev. B 43, 6828 (1991).

¹⁰D. C. Tsui, H. L. Störmer, and A. C. Gossard, Phys. Rev. Lett. 48, 1599 (1982).

¹¹G. Bastard, Appl. Phys. Lett. 43, 591 (1983).

¹²C. Jiang, D. C. Tsui, and G. Weimann, Appl. Phys. Lett. **53**, 1533 (1988).

¹³A. L. Efros, Solid State Commun. 65, 1281 (1988); A. Gold, Appl. Phys. Lett. 54, 2100 (1989).

¹⁴J. A. Davis and J. A. Nixon, Phys. Rev. B 39, 3423 (1989).